

Gravitationally Mediated Entanglement: Newtonian Field vs. Gravitons

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D. Danielson, G.S., & R.M. Wald Phys. Rev. D. 105 086001 (2022) [2112.10798]

D. Danielson, G.S., & R.M. Wald arXiv:2205.06279

D. Danielson, G.S. & R.M. Wald (to appear)

July 4, 2022

Quantum Gravity and Gravitational Entanglement

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What properties of quantum gravity are necessary in order to consistently describe such a quantum superposition?

Gedankenexperiment of Mari *et al.*

- ▶ Two experimenters, Alice and Bob, each control a massive “particle”. There is an **electromagnetic version** (particle is charged) and a **gravitational version** (particle is uncharged).

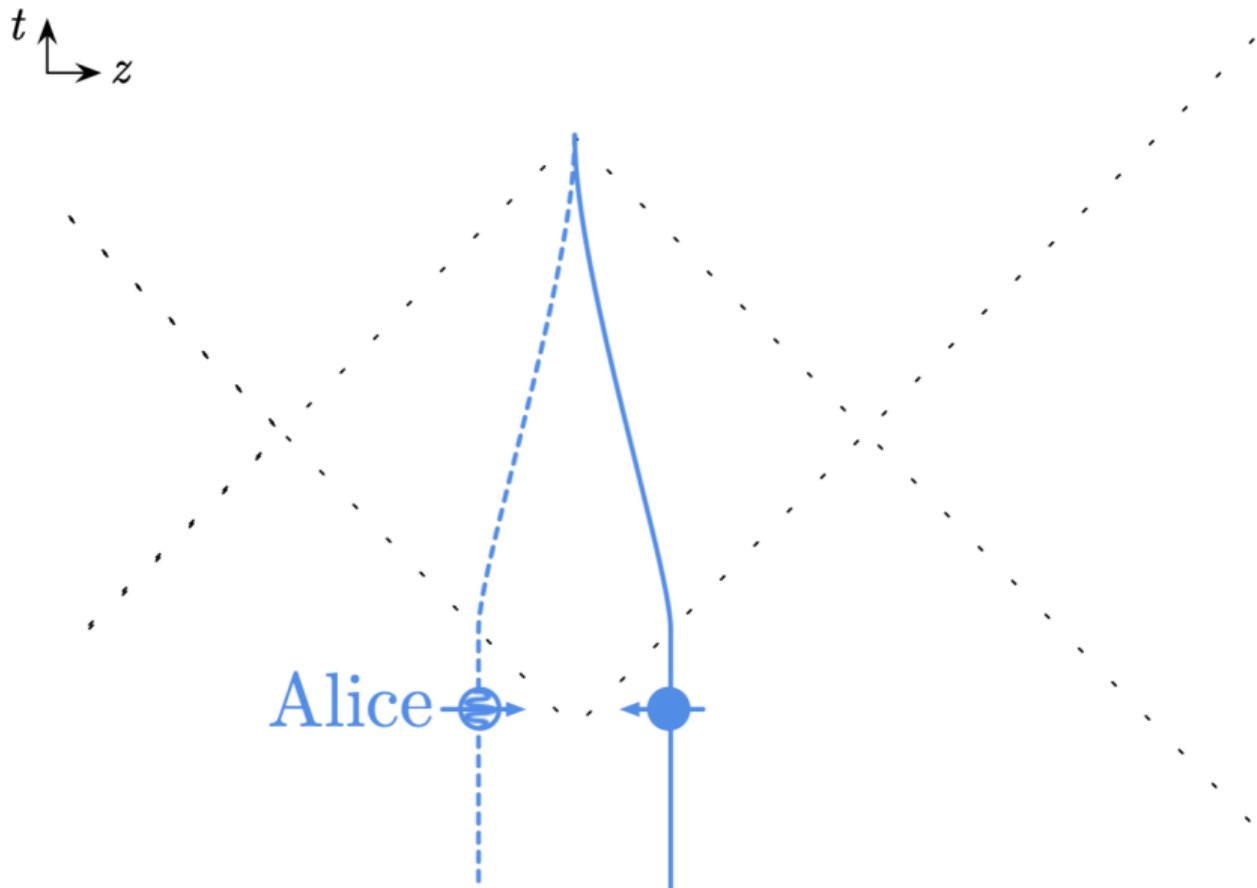
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- ▶ At some (pre-arranged) time, Alice sends her particle through a “reversing Stern-Gerlach apparatus” and then performs an **interference experiment**.

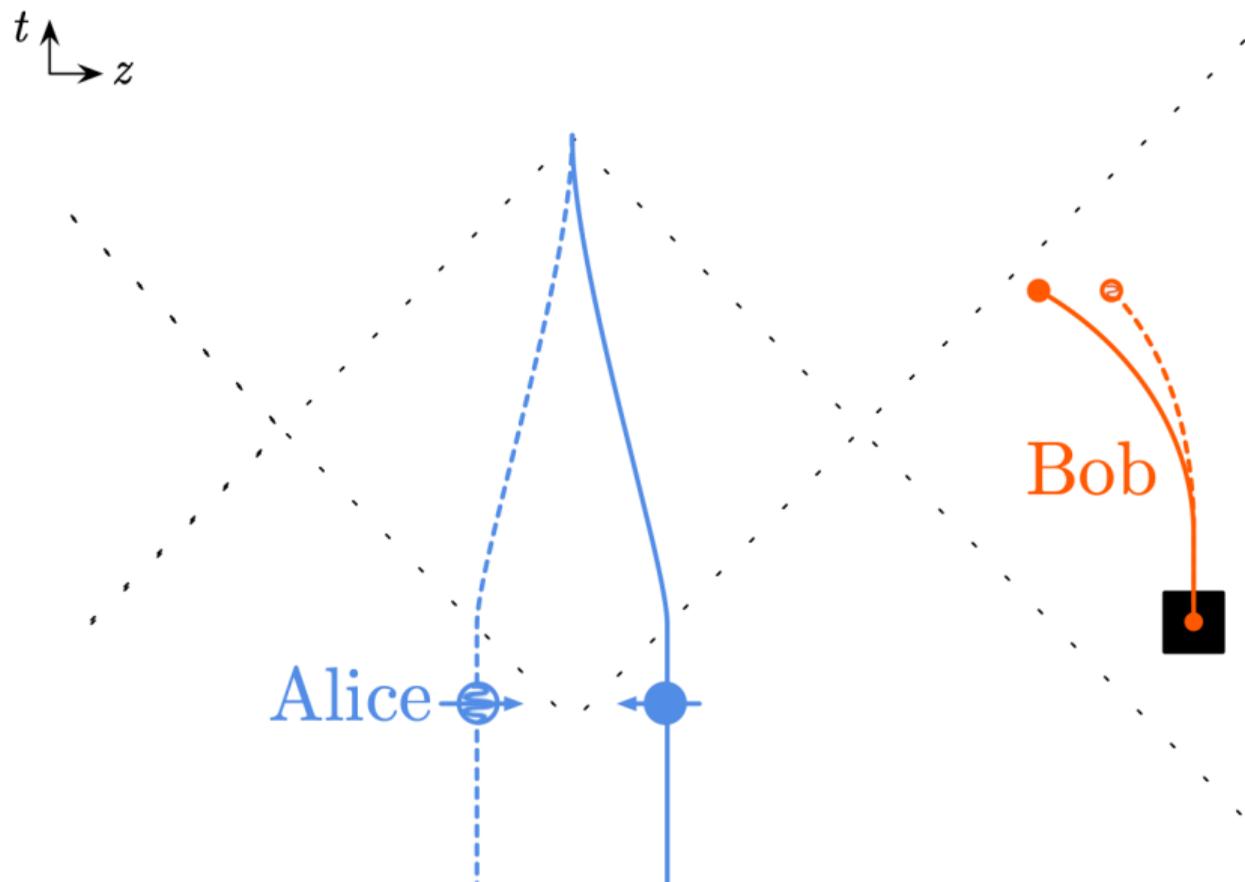
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- ▶ The gravitational version of this gedankenexperiment simply corresponds to replacing "Coulomb" → "Newtonian" and "test charge" → "test mass" ect.

Back-of-the-Envelope Resolution of Belenchia *et al.*

- ▶ Alice must recombine her particles within a light-travel time without emitting any **quantized entangling radiation**. If she accelerates her (charged) superposition too quickly then she will emit an entangling photon/graviton.

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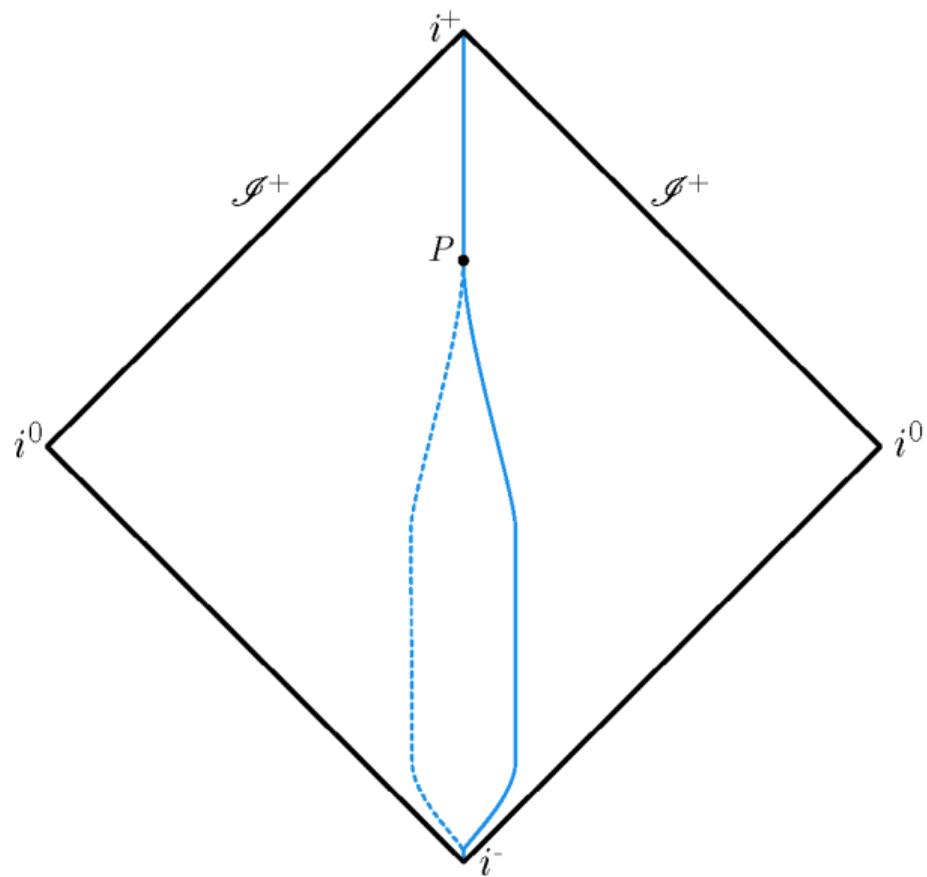
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We now give the precise description of the resolution and gravitational entanglement

Decoherence Due to Alice



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- ▶ Formally, the state of Alice's particle after she makes the spatial superposition is

$$\frac{1}{\sqrt{2}} (|\uparrow; A_1\rangle \otimes |\psi_1\rangle + |\downarrow; A_2\rangle \otimes |\psi_2\rangle)$$

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- ▶ For any time before her recombination, the coherence of Alice's particle is not well-defined. However, at asymptotically late times

$$\frac{1}{\sqrt{2}} (|\uparrow; A_1\rangle_{i^+} \otimes |\Psi_1\rangle_{\mathcal{I}^+} + |\downarrow; A_2\rangle_{i^+} \otimes |\Psi_2\rangle_{\mathcal{I}^+})$$

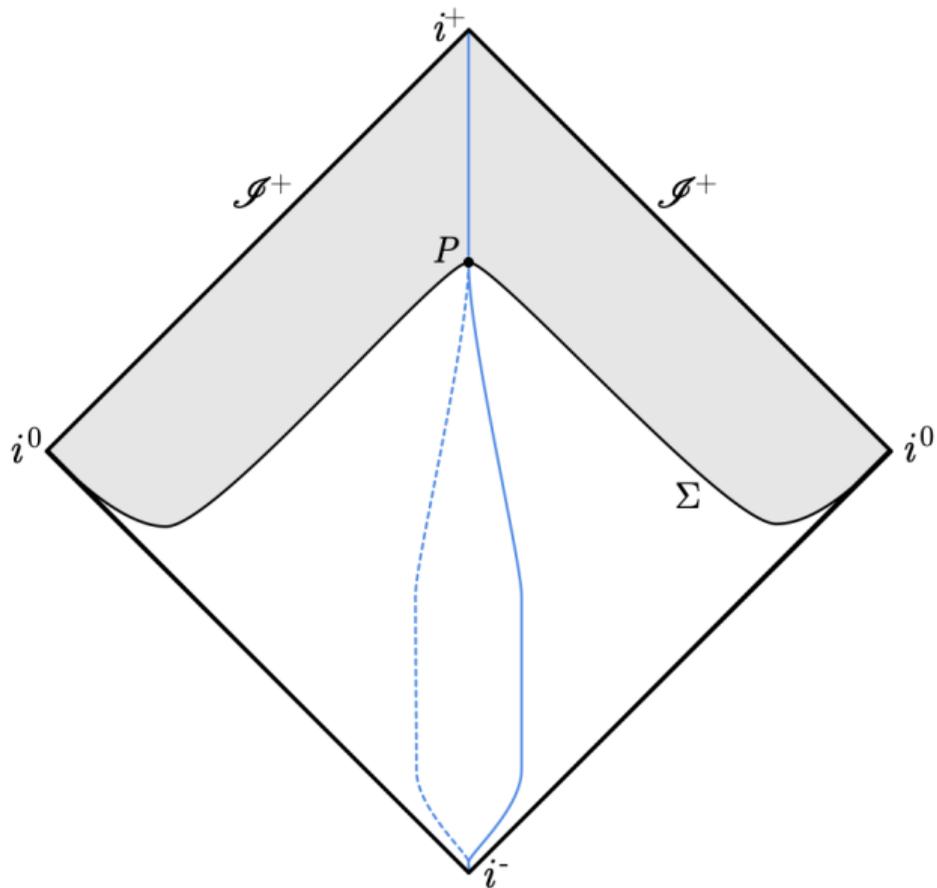
where $|\uparrow; A_1\rangle_{i^+}, |\uparrow; A_2\rangle_{i^+}$ are the states of Alice's particle at i^+ and $|\Psi_1\rangle_{\mathcal{I}^+}, |\Psi_2\rangle_{\mathcal{I}^+}$ are the **free-field, Fock radiation states** at \mathcal{I}^+ for Alice's particle going to the left or right.

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$$\mathcal{D}_{\text{Alice}} = 1 - |\langle \Psi_1 | \Psi_2 \rangle_{\mathcal{H}_+}|$$

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- ▶ After recombining her particle, The state of the field on any Cauchy surface Σ is a common Coulomb field (which does not contribute to her decoherence) and radiation (which is responsible for her decoherence). Propagating the asymptotic states $|\Psi_1\rangle_{\mathcal{I}^+}, |\Psi_2\rangle_{\mathcal{I}^+}$ to the "time" Σ yields the (Coulomb-subtracted) free-field Fock states $|\Psi_1\rangle_{\Sigma}, |\Psi_2\rangle_{\Sigma}$ which cause Alice's decoherence.

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Decoherence Due to Bob

- ▶ Consider the case where Alice would *succeed* in her interference experiment in the absence of external influences (i.e. $|\Psi_1\rangle = |\Psi_2\rangle \approx |0\rangle$). We assume that Bob makes *any* measurement of Alice's field and stops interacting with the field after his measurement.

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$$\frac{1}{\sqrt{2}} (|\uparrow; A_1\rangle_{i+} \otimes |B_1\rangle_{i+} + |\downarrow; A_2\rangle_{i+} \otimes |B_2\rangle_{i+}) \implies \mathcal{D}_{\text{Bob}} = 1 - |\langle B_1 | B_2 \rangle|_{i+}$$

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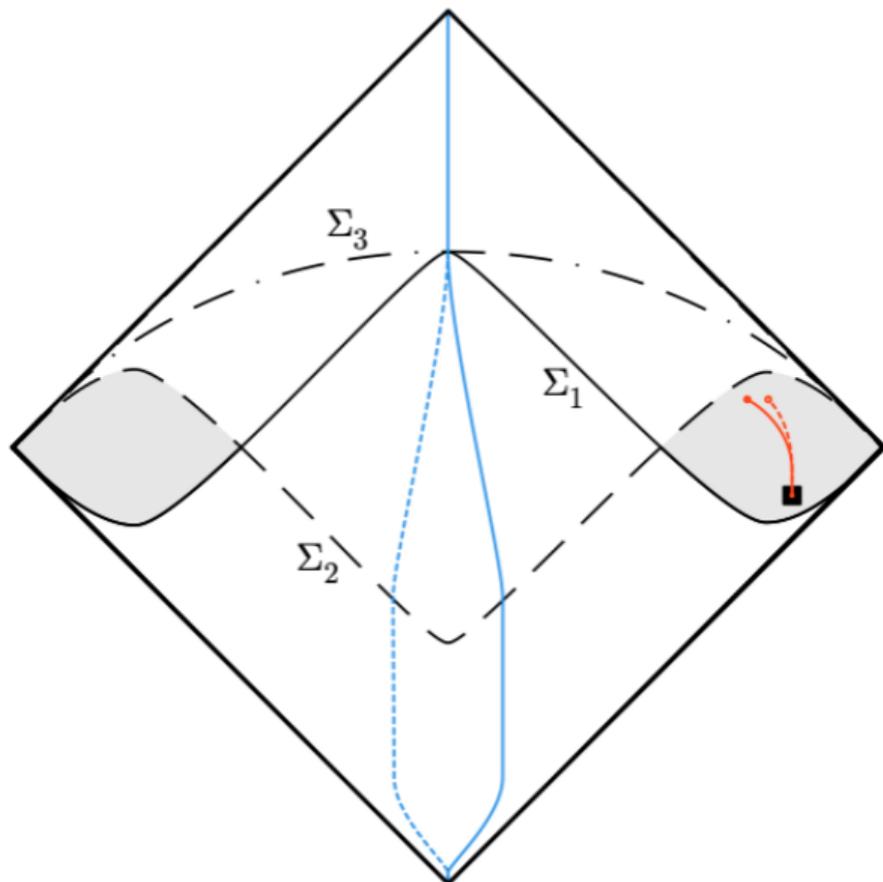
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- ▶ Since Bob completes his measurement after some time T , it follows that $\mathcal{D}_{\text{Bob}} = 1 - |\langle B_1|B_2\rangle|_T$

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- ▶ Consider a Cauchy surface Σ_1 which goes to the **future** of Alice's experiment but to **past** of Bob's experiment

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- ▶ Consider a Cauchy surface Σ_1 which goes to the **future** of Alice's experiment but to **past** of Bob's experiment
- ▶ This decoherence is entirely due to Alice from the (Coulomb-subtracted) **radiation states** emitted during her experiment

$$\mathcal{D}_{\text{Alice}} = 1 - |\langle \Psi_1 | \Psi_2 \rangle|_{\Sigma_1}$$

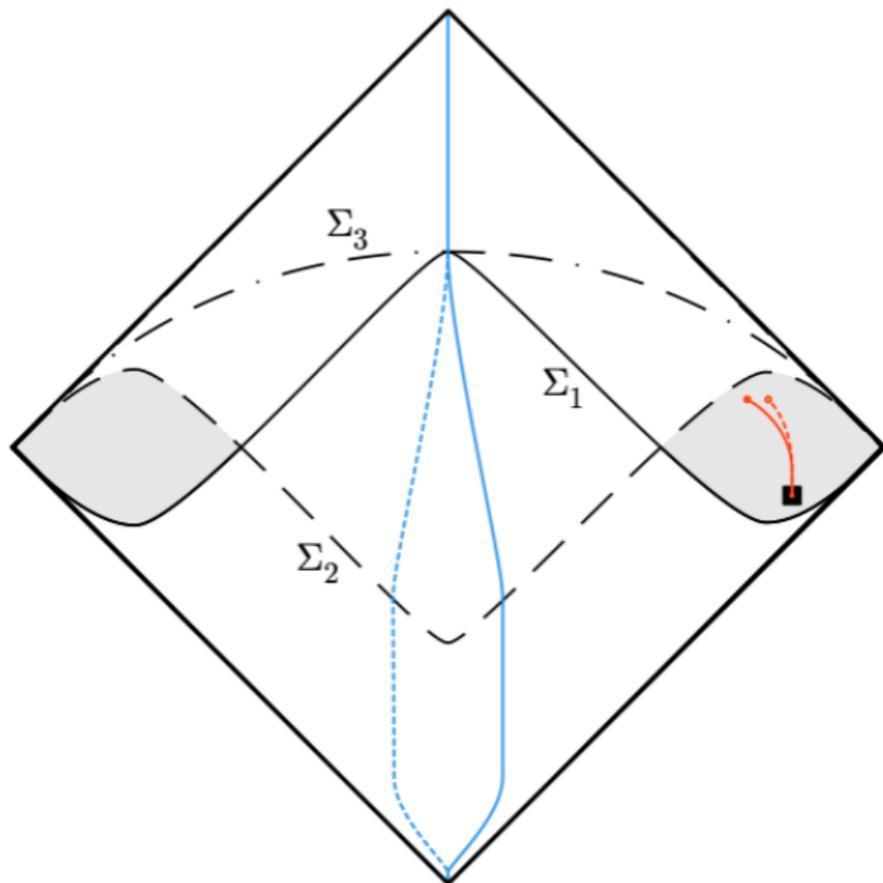
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- ▶ Now consider a surface Σ_2 which goes to the **past** of Alice's experiment but to the **future** of Bob's experiment

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- ▶ Now consider a surface Σ_2 which goes to the **past** of Alice's experiment but to the **future** of Bob's experiment
- ▶ This decoherence is entirely due to the orthogonality of Bob's states from interacting with the stationary Coulomb field of Alice's particle

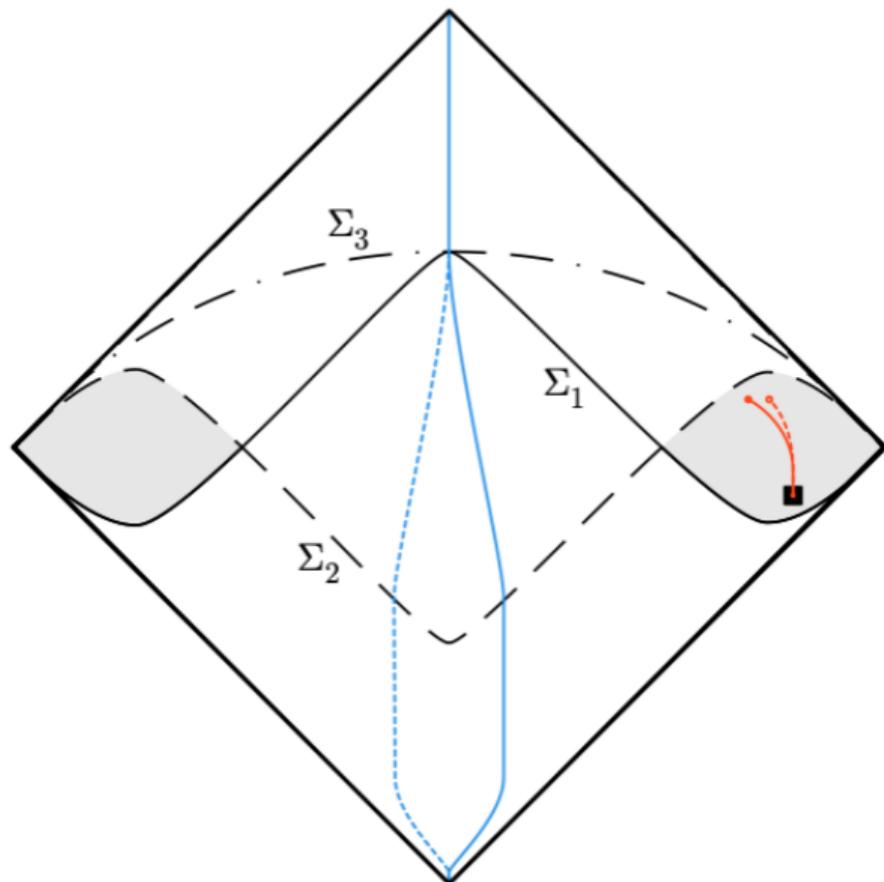
$$\mathcal{D}_{\text{Bob}} = 1 - |\langle B_1 | B_2 \rangle|_{\Sigma_2}$$

There is a paradox if $\mathcal{D}_{\text{Bob}} > \mathcal{D}_{\text{Alice}}$!

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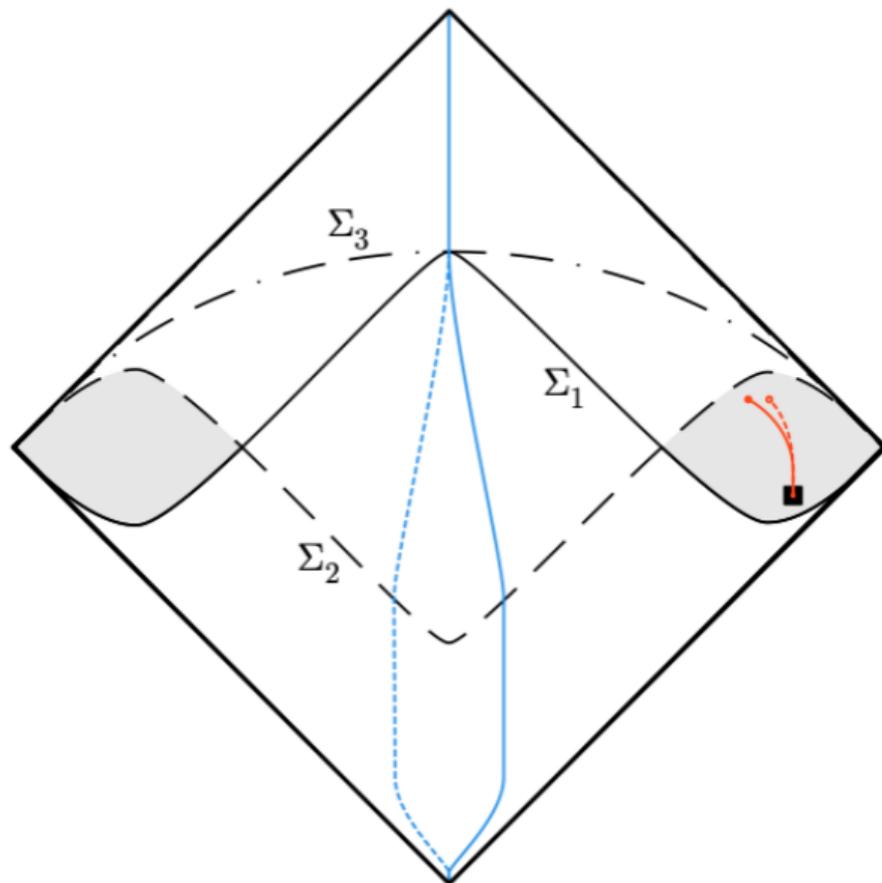
- ▶ This evolution is **unitary**: $\langle \Psi'_1 | \Psi'_2 \rangle_{\Sigma_3} \langle B_1 | B_2 \rangle = \langle \Psi_1 | \Psi_2 \rangle_{\Sigma_1} \langle B_0 | B_0 \rangle = \langle \Psi_1 | \Psi_2 \rangle_{\Sigma_1}$

$$|\langle B_1 | B_2 \rangle| \geq |\langle \Psi_1 | \Psi_2 \rangle_{\Sigma_1}| \implies \mathcal{D}_{\text{Bob}} \leq \mathcal{D}_{\text{Alice}}$$

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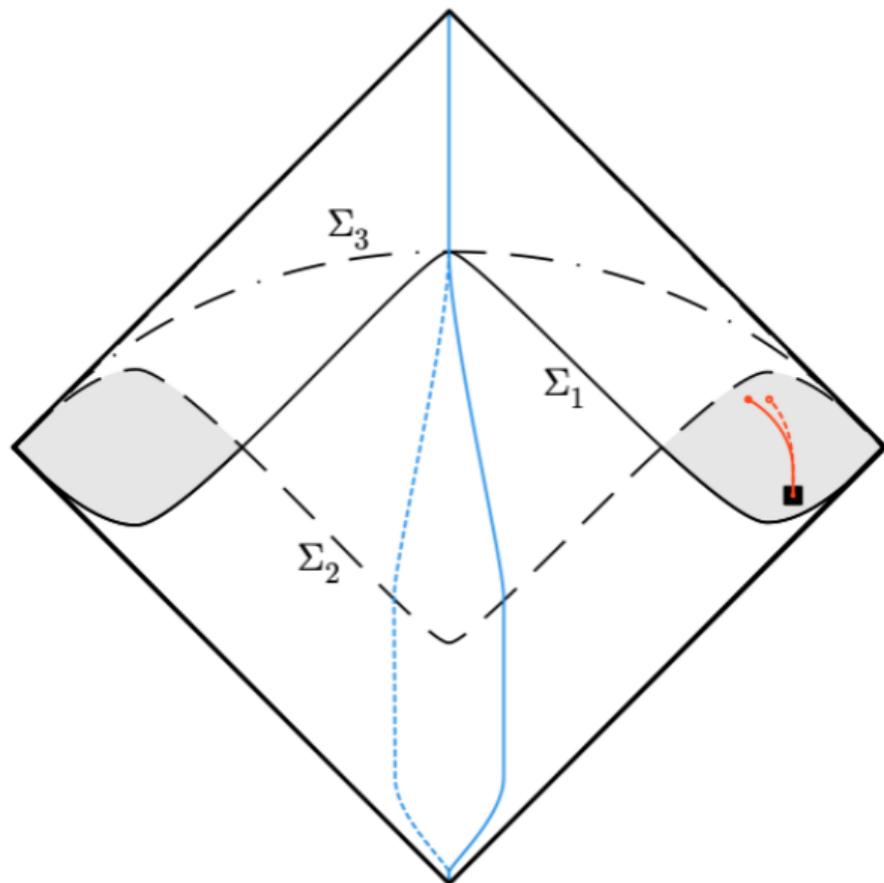
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- ▶ For spacelike separated measurements, there is no clear distinction between entanglement via the Newtonian field and entanglement via on-shell gravitons. *Both viewpoints are equivalent! Both viewpoints must be valid descriptions of the process in order to have a consistent description of a quantum superposition.*
- ▶ This suggests that Newtonian entanglement implies the existence of graviton entanglement and *the experimental discovery of Newtonian entanglement can be viewed as implying the existence of the graviton*

- ▶ Gedankenexperiments where both quantum theory and gravity play an essential role can give considerable insight into the nature of quantum gravity.
- ▶ Many actual experiments have been recently proposed to measure entanglement mediated by the Newtonian field. It has been claimed by a number of authors that such experiments cannot teach us anything about quantum gravity since it doesn't couple to the “true propagating degrees of freedom”.
- ▶ Our re-analysis of the gedankenexperiment by Mari et al shows that the consistent description a quantum superposition shows that **Newtonian entanglement implies the existence of graviton entanglement**. Thus, this strongly supports the view that the success of recently posed experiments implies the existence of the graviton as a fundamental particle.