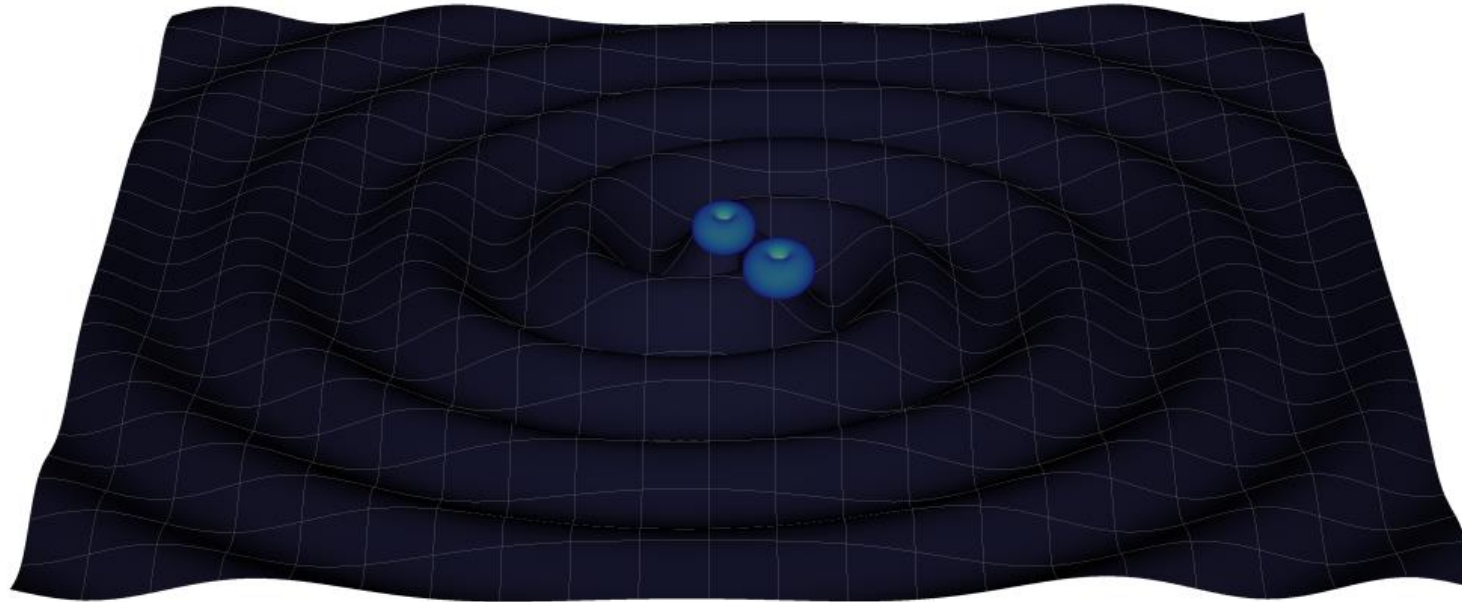


Constraining fundamental properties of boson stars through their multipole moments



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Motivations for the work

- Gravitational waves allow to probe the nature of compact objects and to search for new physics

*Exotic
Compact
Objects*



Boson stars



Proca stars



BHs with scalar hair



Fuzzballs . . .

- The multipolar structure affects the dynamics of binary systems and their gravitational wave emission

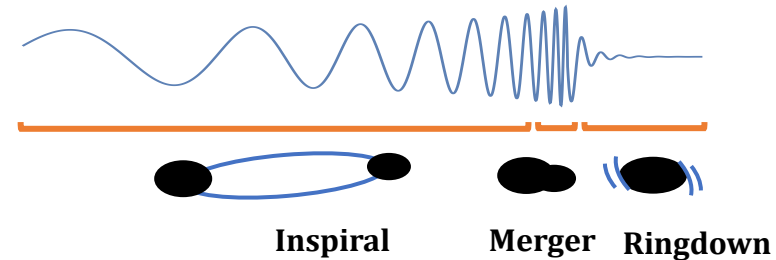
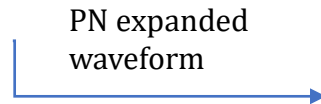
Stationary axisymmetric spacetime \Rightarrow scalar mass moments M_0, M_2, \dots and current moments S_1, S_3, \dots

for a Kerr black hole

$$M_l + iS_l = M^{l+1} (i\chi)^l \quad \chi = \frac{J}{M^2}, \quad M = M_0$$

Not true for a generic compact object!

E.g. $h \sim \mathcal{A}(f) e^{i(\psi_{BH}(f) + \text{finite size corr})}$



Rotating Boson Stars (BSs)

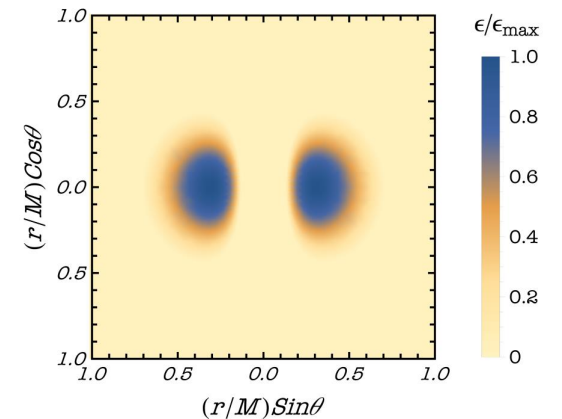
- Boson stars are solutions of the Einstein gravity, minimally coupled to a complex scalar field:

$$L = -\frac{1}{2}g^{\mu\nu}\phi_{,\mu}\phi_{,\nu} - \frac{1}{2}V(|\phi|^2)$$

The equations take the form: $G_{ab} = 8\pi T_{ab}$ $\frac{1}{\sqrt{-g}}\partial_a(\sqrt{-g}g^{ab}\partial_b\phi) = \frac{dV(|\phi|^2)}{d|\phi|^2}\phi$

- Neutron Stars: Equation Of State* → *Boson Stars: Self-interactions in $V(|\phi|^2)$*

- Mini BSs	$V(\phi ^2) = m^2 \phi ^2$	$M_{max} \sim \frac{M_p^2}{m}$
- Massive BSs	$V(\phi ^2) = m^2 \phi ^2 + \lambda \phi ^4$	$M_{max} \sim \frac{M_p^3}{m^2}\lambda^{\frac{1}{2}}$
- Solitonic BSs	$V(\phi ^2) = m^2 \phi ^2 \left(1 - \frac{2 \phi ^2}{\sigma^2}\right)$	$M_{max} \sim \frac{M_p^4}{m\sigma^2}$



Normalized energy-density of a BS in a transversal section

Theoretical set up

- Our case is: $V(|\phi|^2) = \frac{1}{2}m^2|\phi|^2 + \frac{1}{4}\lambda|\phi|^4$ in the large self-coupling limit ($\lambda \gg m^2$)
- In the $\lambda \gg m^2$ a reasonable assumption is : $\phi_{,r} \sim 0 \quad \phi_{,\theta} \sim 0 \quad \phi_{,\varphi} \neq 0 \Rightarrow$

1) The stress energy tensor assume the form:

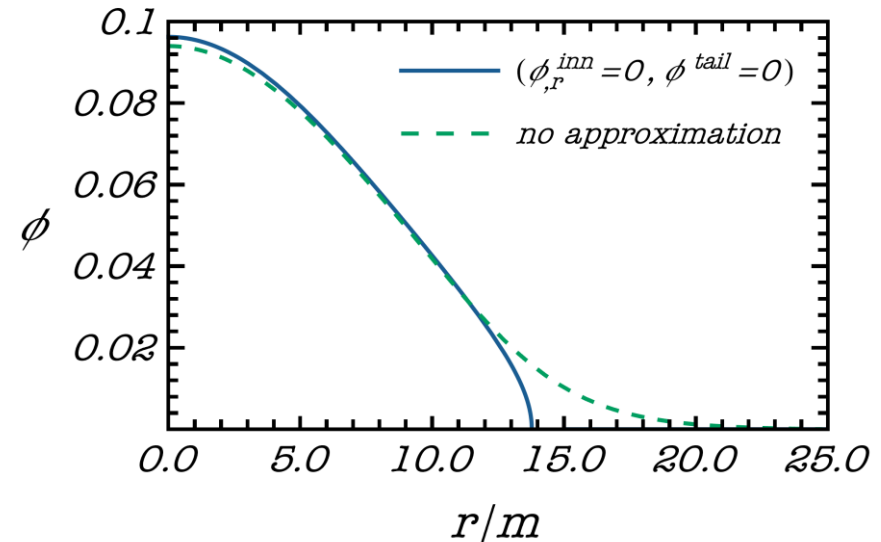
$$T_{\mu\nu} = (\epsilon + P)u_\mu u_\nu + P g_{\mu\nu}$$

2) The scalar field eq. becomes trivial:

$$|\phi|^2 = \text{Max}[0, \lambda^{-1}(-g^{tt}\Omega^2 + 2g^{t\varphi}\Omega s - g^{\varphi\varphi} - m^2)]$$

3) The coupling constants can be removed:

$$\frac{1}{\lambda^{\frac{1}{2}}/m^2} \equiv M_B$$



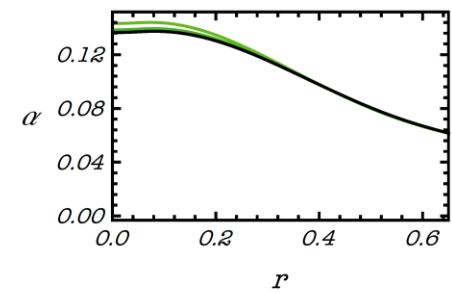
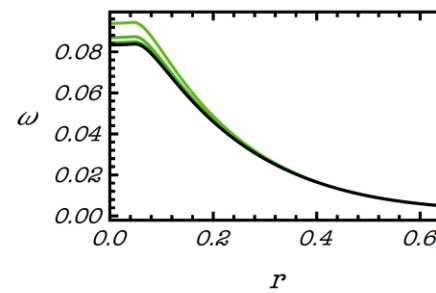
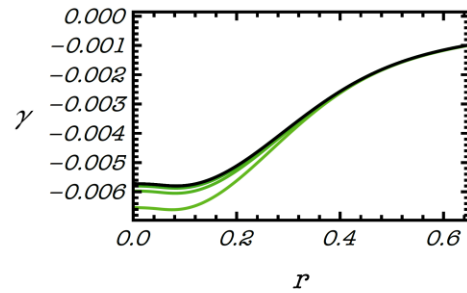
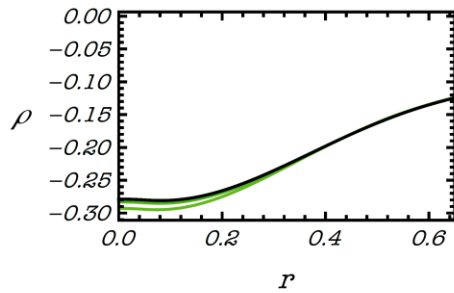
In the figure a similar approximation is considered for the non-spinning case.

Self-consistent field method

- In these conditions, one can apply a self-consistent field method to BSs.

F. D. Ryan, Phys. Rev. D 55, 6081 (1997), Spinning boson stars with large self-interaction.

- $ds^2 = -e^{\gamma(r,\theta)+\rho(r,\theta)} dt^2 + e^{2\alpha(r,\theta)} (dr^2 + r^2 d\theta^2) + e^{\gamma(r,\theta)-\rho(r,\theta)} r^2 \sin^2 \theta (d\varphi - \omega(r,\theta) dt)^2$



Cycle — 10 — 20 — 30 — 40 — 50 — 60 — 70 — 80 — 90 — 100 — 110 — 120 — 130 — 140 — 150

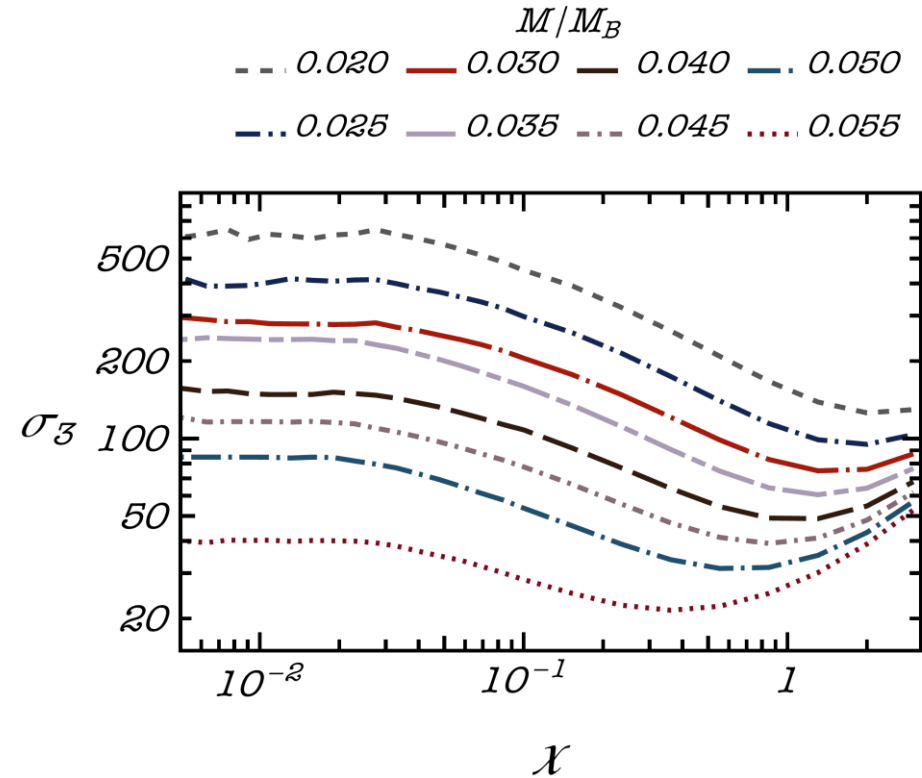
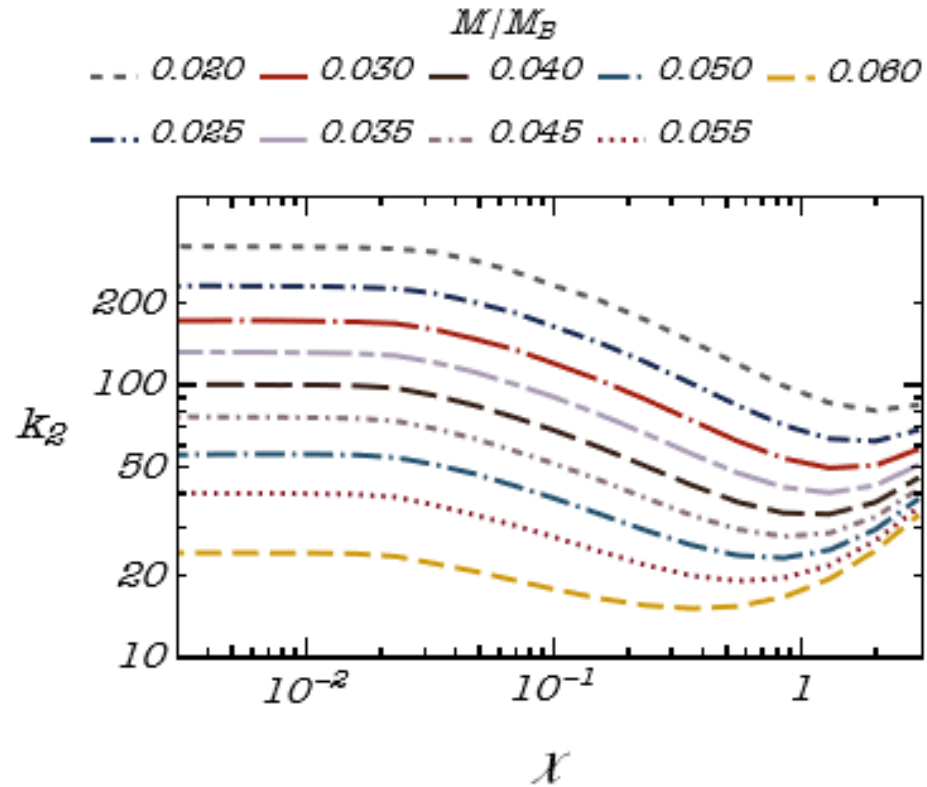
- Mass and current moments $\{M_0, M_2, \dots\}$, $\{S_1, S_3, \dots\}$ can be read off:

$$\rho(r, \mu) = \sum_{n=0}^{\infty} -2 \frac{M_{2n}}{r^{2n+1}} P_{2n}(\mu) + \text{higher orders}$$

$$\omega(r, \mu) = \sum_{n=1}^{\infty} -\frac{2}{2n-1} \frac{S_{2n-1}}{r^{2n+1}} \frac{P_{2n-1}^1(\mu)}{\sin \theta} + \text{higher orders}$$

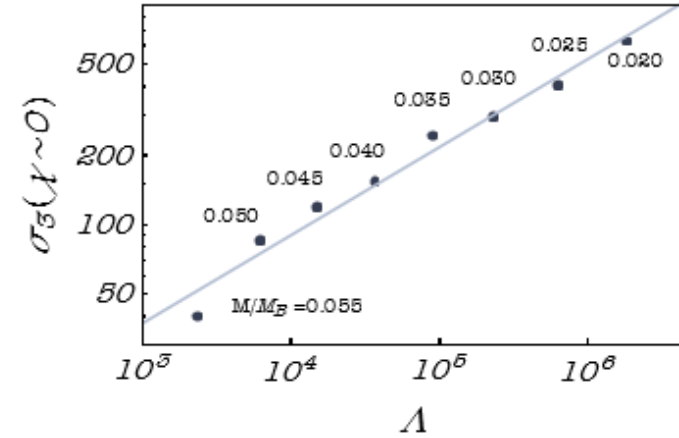
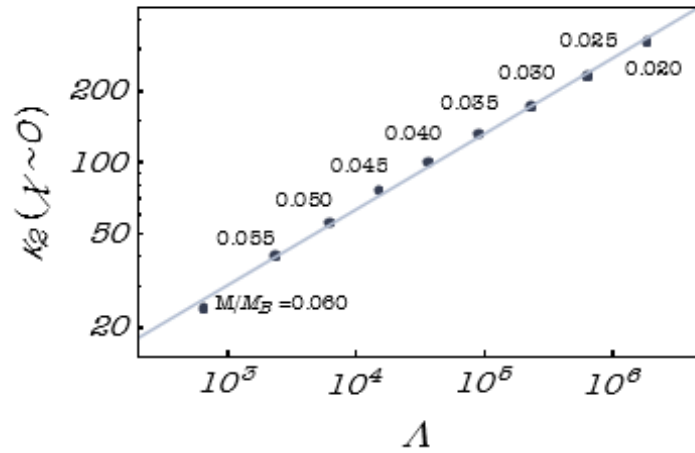
Mass-quadrupole and spin-octupole moment

- The plot shows $\kappa_2 = -\frac{M_2}{\chi^2 M^3}$ and $\sigma_3 = -\frac{S_3}{\chi^3 M^4}$ as a function of the dimensionless spin χ for different BS masses (for a Kerr BH $\kappa_2 = \sigma_3 = 1$).

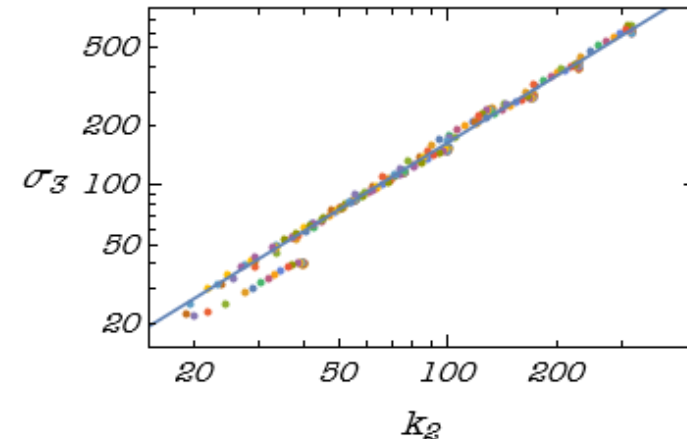
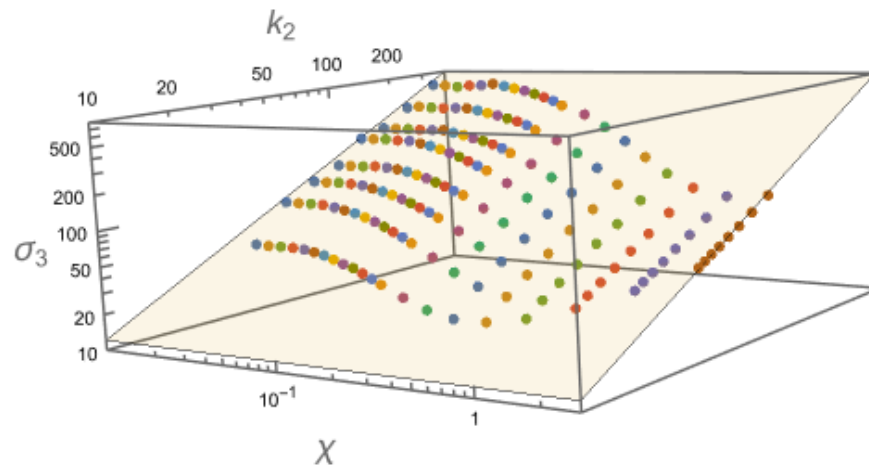


Universal Relations for Boson Stars?

- The reduced quadrupole and octupole moments are simply connected to the tidal deformability



- The relation between κ_2 and σ_3 appears remarkably to be independent on the spin χ



Conclusions and perspectives

We have analyzed the characteristics and the multipolar structure of rotating BSs with quartic self-interactions in the strong coupling limit ($\lambda \gg m^2$):

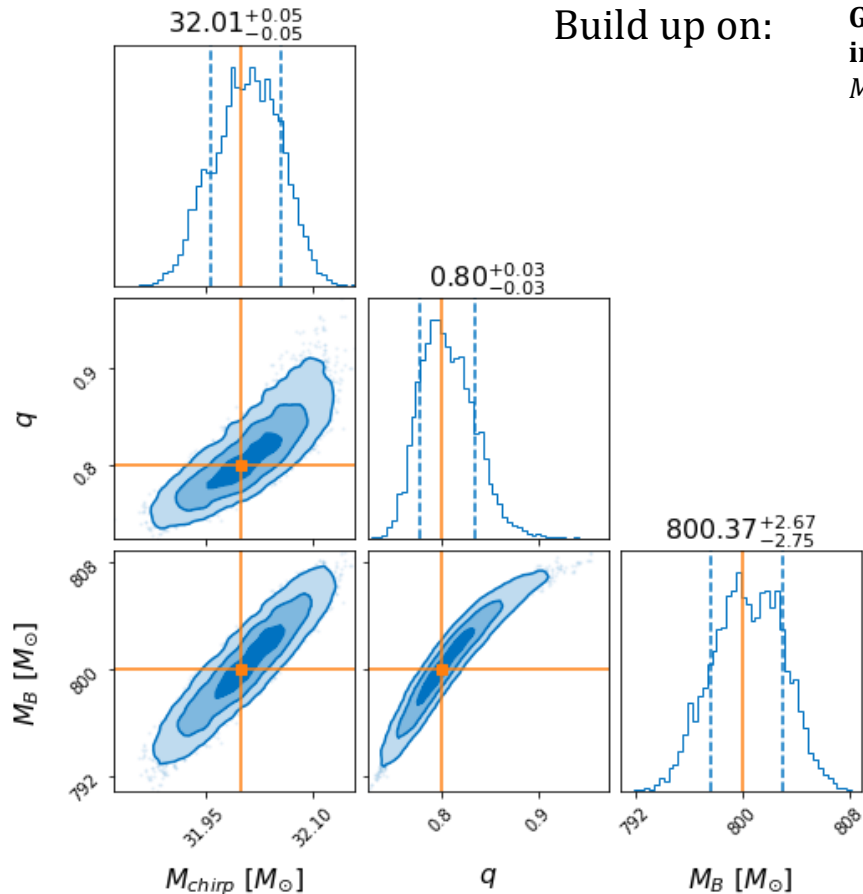
- The values of the quadrupole and octupole moments agree with previous results for high spins. In the low spin region, we confirmed the relations $M_2 \sim \chi^2$ ($S_3 \sim \chi^3$), but with improved accuracy, we found absolute values which are larger by a factor of ~ 2 (~ 3).
- We found that simple relations hold linking the multipole moments and the tidal deformability.

Future perspectives and extensions of this work include:

- Study the impact of different self-interactions on the multipolar structure to investigate:
 - the universality degree of the relations between multipole moments and tidal deformability
 - the possibility of distinguish the gravitational signal from different BS source models.

Ongoing work

- Build a coherent waveform model which consistently includes different contributions of the multipole moments and assess their detectability through current/future GW observations



Gravitational-wave detectors as particle-physics laboratories: Constraining scalar interactions with a coherent inspiral model of boson star binaries, *Costantino Pacilio, Massimo Vaglio, Andrea Maselli, Paolo Pani. Phys.Rev. D 102 (2020) 8, 083002*

TaylorF2 coherent waveform with:

$$Q = -\kappa(\chi, M/M_B)\chi^2 M^3$$

and Λ from:

$$\frac{M}{M_B} = \frac{\sqrt{2}}{8\sqrt{\pi}} \left[-0.828 + \frac{20.99}{\log \Lambda} - \frac{99.1}{(\log \Lambda)^2} + \frac{149.7}{(\log \Lambda)^3} \right]$$

N. Sennett et al., Phys. Rev. D, 96, 2 (2017) 024002

Tidal deformability and quadrupoles are functions of mass, spin and M_B !