

INFRARED FINITE SCATTERING IN QFT & QUANTUM GRAVITY

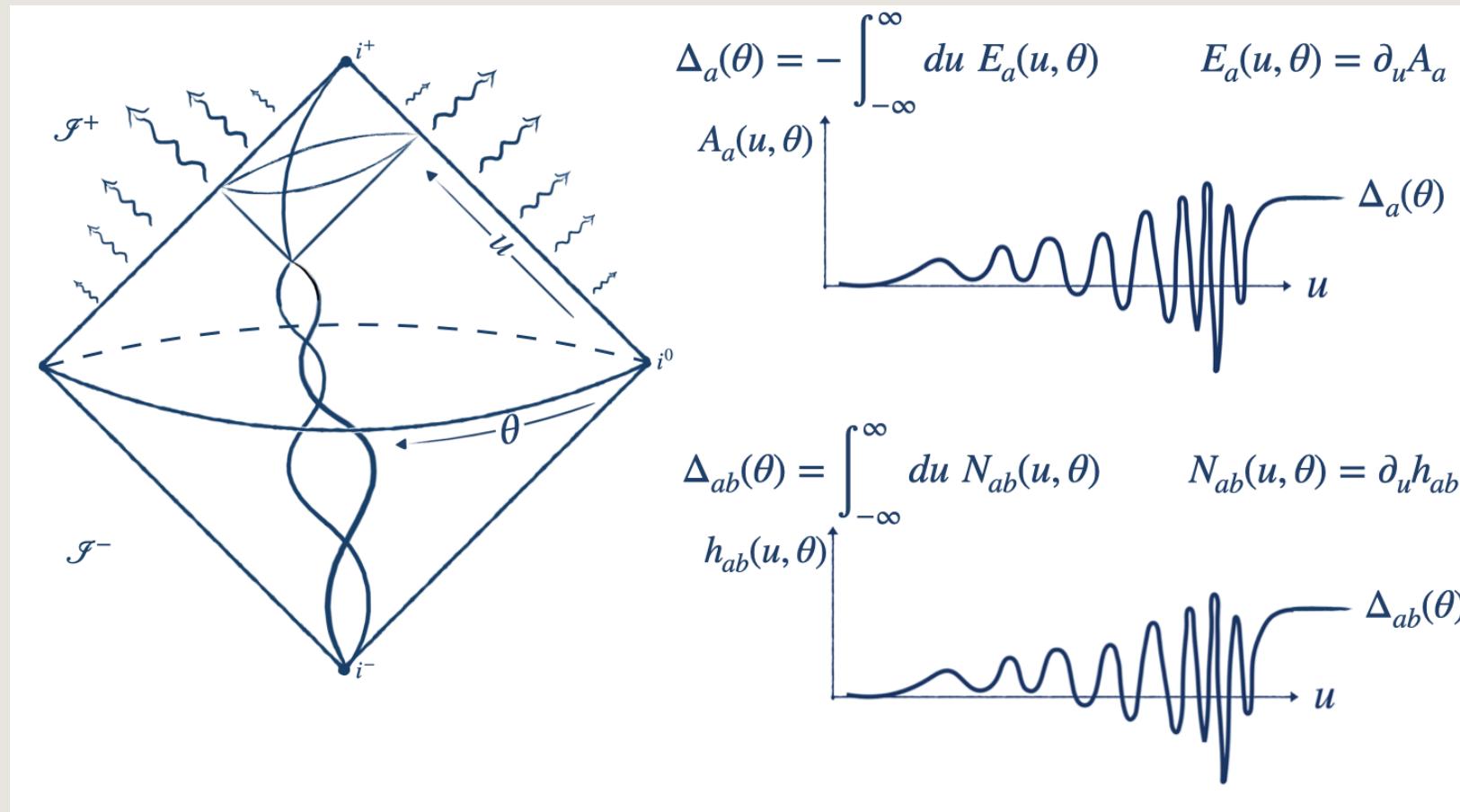
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SCATTERING: ASYMPTOTIC FLATNESS

1. generic scattering of particles or fields from past to future infinity
2. The radiation field has a long time tail or a zero frequency part



MASSIVE QED: ASYMPTOTIC FIELDS

Massive charged scalar field φ coupled to electromagnetism

- Massive scalar asymptotics in the limit to i^\pm described by $b(p), b^\dagger(p), c(p), c^\dagger(p)$, where p lies on a unit hyperboloid. (*usual Fourier description*)
- EM vector potential near \mathcal{I}^\pm :
 $A_a \sim A_A(u, \theta) + O(1/r)$
Radiative field: $E_A = \partial_u A_A$
- EM memory:
 $\Delta_A(\theta) \sim A_A(u \rightarrow +\infty) - A_A(u \rightarrow -\infty) = - \int du E_A(u, \theta)$

EM: FOCK SPACE

Quantum operator-valued *distribution*: $\mathbf{E}_A(u, \theta)$.

- Vacuum state: $|\omega_0\rangle$ (*Poincaré invariant*): $\langle\omega_0|\mathbf{E}_{1A}\mathbf{E}_{2B}|\omega_0\rangle \propto \frac{\delta(\theta_1, \theta_2)}{(u_1 - u_2 - i0^+)^2} q_{AB}$
- one-photon Hilbert space: test “wavefunctions” $s_A(u, \theta)$

$$\|s\|^2 \propto \int d^2\theta \int_0^\infty d\omega \omega |s(\omega, \theta)|^2$$

- If wavefunction has memory then $s_A \sim \frac{\Delta_A(\theta)}{\omega}$ and $\|s\|^2$ *diverges logarithmically*.
No states with memory in standard Fock space!

EM: MEMORY HILBERT SPACES

Use coherent states with memory

- Classical solution $\mathcal{E}_A(u, \theta)$ with memory $\Delta_A(\theta) = - \int du \mathcal{E}_A(u, \theta)$
- $\mathbf{E}_A \mapsto \mathbf{E}_A + \mathcal{E}_A \mathbf{1}$ (*automorphism of the operator algebra*)
- New Hilbert space of states \mathcal{F}_Δ with memory $\Delta_A(\theta)$
- \mathcal{F}_Δ is *unitarily-equivalent* to $\mathcal{F}_{\Delta'}$ if and only if $\Delta_A(\theta) = \Delta'_A(\theta)$
- *Uncountably-many* inequivalent Hilbert spaces!

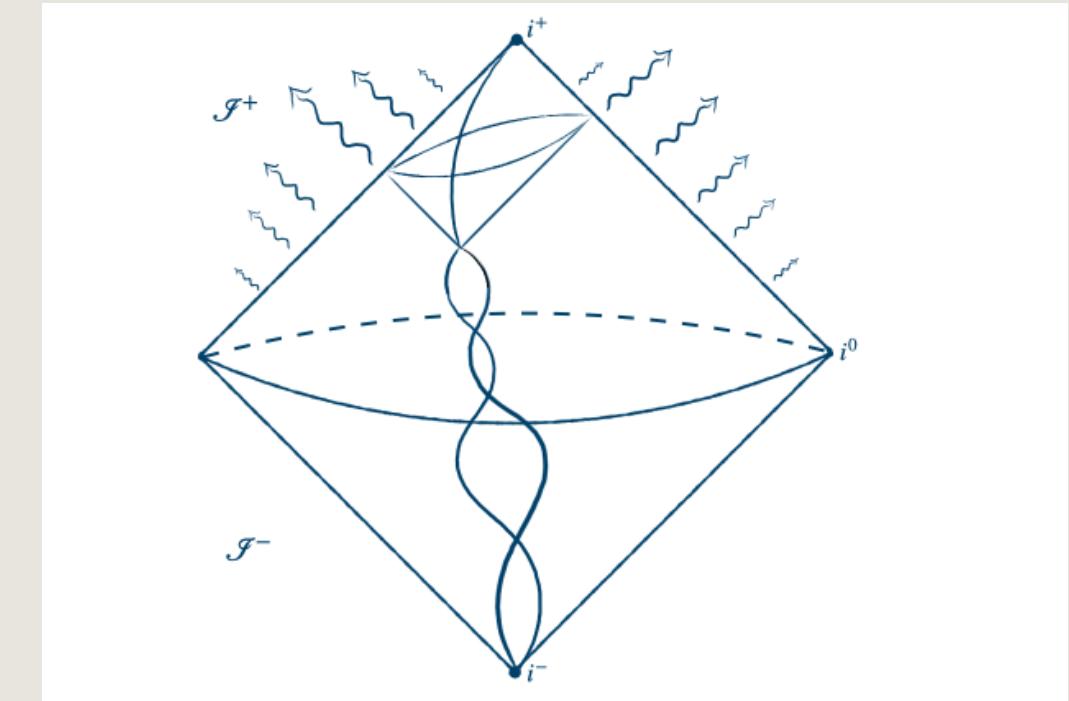
MASSIVE QED: SCATTERING

- EM with no sources: $\Delta_A^{out}(\theta) = \Delta_A^{in}(-\theta)$; antipodal map $\theta \mapsto -\theta$ on \mathbb{S}^2
- EM with *classical* source: $\Delta_A^{out}(\theta) = \Delta_A^{in}(-\theta) + \text{classical current term}$
- EM with *massive quantum* field: given a $\Delta_A^{in}(\theta)$ *no* definite out-memory!
- If we *pretend* that all states are in the standard Fock space \mathcal{F}_0 then S-matrix diverges (*Weinberg soft theorem*)

Idea: build Hilbert spaces using *conserved* quantities, not memory.

CONSERVED CHARGES AT SPATIAL INFINITY

- $\mathcal{Q}_{i^0}(\theta) = \mathcal{Q}_{i^-}(\theta) + \mathcal{D}^A \Delta_A(\theta)$
- $\mathcal{Q}_{i^0}^{out}(\theta) = \mathcal{Q}_{i^0}^{in}(-\theta)$



Find Hilbert spaces of definite charge $\mathcal{Q}_{i^0}(\theta)$

FADDEEV-KULISH HILBERT SPACE

1. add *any* EM state $|\psi\rangle \in \mathcal{F}_\Delta$ with memory chosen so that

$$\mathcal{Q}_{i^0}(\theta) = 0 = \mathcal{Q}_{i^-}(\theta; \{p_i, q_i\}) + \mathcal{D}^A \Delta_A(\theta; \{p_i, q_i\})$$

2. Then the *dressed* state

$\int d^3 p_1 \dots d^3 q_n w(\{p_i, q_i\}) |p_1, \dots, p_n, q_1, \dots, q_n\rangle \otimes |\psi\rangle_{\Delta(\theta; \{p_i, q_i\})}$ has zero charge at spatial infinity.

3. Faddeev-Kulish Hilbert space \mathcal{H}^{FK} by direct summing over particle-antiparticle number.

Charge conservation ensures that this Hilbert space scatters to itself! So the S-matrix $S : \mathcal{H}^{FK} \rightarrow \mathcal{H}^{FK}$ is well-defined.

FK HILBERT SPACES: ISSUES

- Total charge cannot be dressed away; memory is $\ell \geq 1$; always need equal number of particles and antiparticles (*hide particles behind the moon*)
- Can define FK Hilbert spaces with any value of $Q_{i^0}(\theta)$ but not Lorentz-invariant; angular momentum/spin of states is undefined! (“*Lorentz is spontaneously broken*”)
- Massive fields always come with a radiative photon cloud; can hide the cloud at very low frequencies but always have to send it in.

MASSLESS QED & LINEARIZED GRAVITY

- Massless QED: solve $\mathcal{Q}_{i^0}(\theta) = 0 = \mathcal{J}(\theta; \{p_i, q_i\}) + \mathcal{D}^A \Delta_A(\theta; \{p_i, q_i\})$
- $\mathcal{J}(\theta; p) = \delta(\theta, \hat{p})$ implies $\Delta_A \sim 1/(\theta - \hat{p})$ (*collinear divergence*); Δ_A is not L^2 ; *photon dressing has infinite energy!*
- Linearized gravity: solve $\mathcal{Q}_{i^0}(\theta)|_{\ell \geq 2} = 0 = \mathcal{T}(\theta; \{p_i, q_i\}) + \mathcal{D}^A \mathcal{D}^B \Delta_{AB}(\theta; \{p_i, q_i\})$
- $\mathcal{T}(\theta; p) = \delta(\theta, \hat{p})$ implies $\Delta_{AB} \sim \log(\theta - \hat{p})$ (*collinear divergence*); Δ_{AB} is L^2 ; *dressing has finite energy!*
- $\mathcal{Q}_{i^0}(\theta)|_{\ell \geq 2} = 0$ is not Lorentz-invariant, so no angular momentum/spin for these states

NONLINEAR GR

Radiative field is News N_{AB} ; gravitational memory is $\Delta_{AB}(\theta) = \int du N_{AB}(u, \theta)$

$$Q_{i^0}(\theta) = \int du N_{AB}(u, \theta) N^{AB}(u, \theta) + \mathcal{D}^A \mathcal{D}^B \Delta_{AB}(\theta)$$

- The memory Hilbert spaces \mathcal{F}_Δ still exist and memory is not conserved
- *gravity sources itself*; any kind of dressing will further contribute to the charge

Theorem: The only *Hadamard* state of fixed $Q_{i^0}(\theta)$ is the *vacuum* with $Q_{i^0}(\theta) = 0$.

So no non-vacuum state is in *any* FK Hilbert space!

ALGEBRAIC SCATTERING

- In/out states defined as lists of correlation functions $\omega(\mathbf{O}_1, \dots, \mathbf{O}_n)$
- Any such state lives in some Hilbert space, but *not all states* can be put in a *single Hilbert space*
- Evolution/scattering is given by a \$-map: $\$: \omega^{in} \mapsto \omega^{out}$
- conservation of probability: $\omega^{in}(1) = 1$ implies $\omega^{out}(1) = 1$;
no information loss: $\$$ maps pure in-states to pure out-states

How to *practically* compute such a $\$$?  Usual methods like Feynman diagrams/LSZ reduction do not work; Hawking (1976) has interesting ideas, but still assumes a single Hilbert space

don't take a vacation yet! Or maybe take a vacation and think about it!!