

A novel formulation for the evolution of relativistic rotating stars



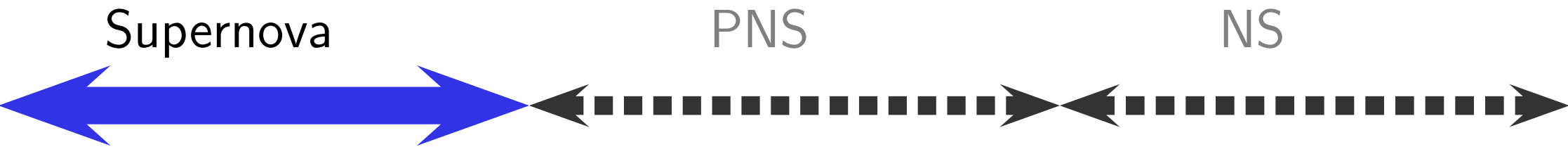
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Based on arXiv[gr-qc]: 2204.09943

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Supernova explosion



Core-collapse supernovae ($M \geq 8M_{sun}$)

Many References...

(Marek&Janka '09, Müller+ '12, Bruenn+ '13,'16 Takiwaki+ '14, Dolence+ '15, Lentz+ '15', Melson+ '15, O'Connor&Cough '15, Just+ '15, Kuroda+ '16, Summa+ '16, Pan+ '16, Roberts+ '16, Burrows+ '16, Andresen+ '17, ...)

Neutrino Transport:

First-principle calculation by Boltzmann equation

Successful explosion:

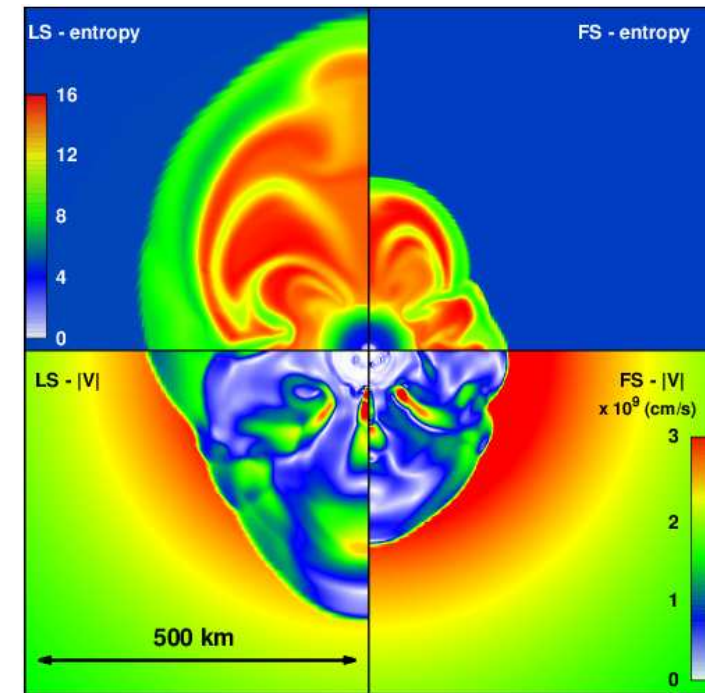
Progenitor, EOS, Rotation, Dimension, GR?

Harada+ '19&'20, Akaho+ '21, Iwakami+ '21, Akaho+ '22

1 model (axisymmetry) \sim 6 month by supercomputer

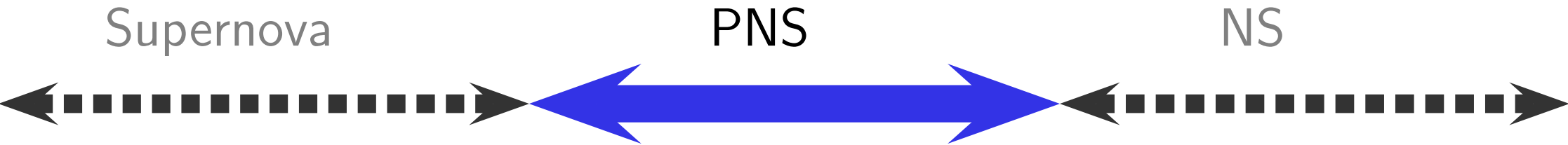
\rightarrow 3D simulations in Fugaku Iwakami+ '20

- Dynamical simulation (~ 10 [sec]) $R \sim 50$ [km]
- Neutrino Cooling (Diffusion) Timescale $\sim \mathcal{O}(10)$ [sec]



Nagakura *et al.* (2018)

Evolution of Proto-Neutron Star(PNS)



○ Lagrange Evolution of PNS Burrows&Lattimer '86, Keil+ '96, Pons+ '99

- Neutrino Transfer
→ Variation of Electron fraction
- Energy Equation
→ Variation of Entropy
- Equilibrium
(Relativistic star with spherical symmetry)

○ Euler “Evolution” of Rotating Stars in GR

Goussard+ '97&'98, Sumiyoshi+ '99, Strobel+ '99, Villain+ '04

- Rotation Law: Uniform or Differential Rotation ($+\alpha$)

Equilibrium of Rotating Stars in GR

Hartle '67
 Butterworth&Ipsen '76
 Komatsu+ '89
 Gourgoulhon+ '99, ...
 Uryu+ '17, Camelió+ '19

○ A perfect fluid in GR (Force-Balance eq.)

○ Geometry: Second-order PDEs for $g_{\mu\nu}$

$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

○ Fluid: First-order PDEs for P (or ρ)

$$\frac{1}{\varepsilon + P} \nabla_{\mu} P = \nabla_{\mu} \log(u^t) - F \nabla_{\mu} \Omega$$

Pressure Gradient

Gravity

Centrifugal Force

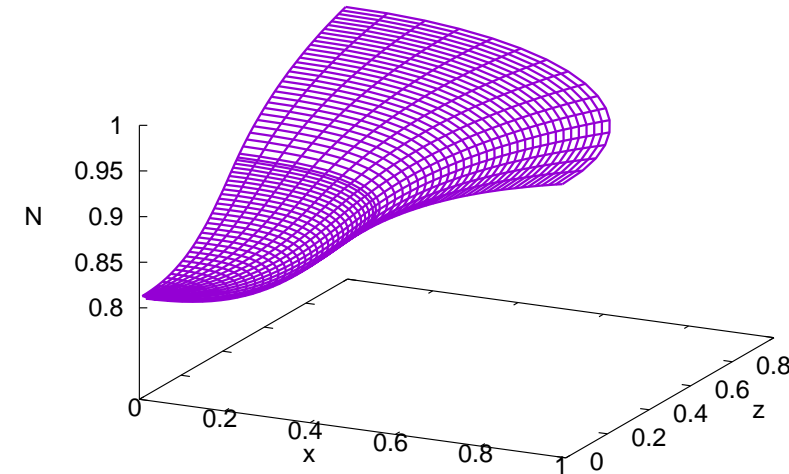
EOS: (pseudo-)Barotrope

$$H(p) \equiv \int_0^P \frac{dP'}{\varepsilon(P') + P'}$$

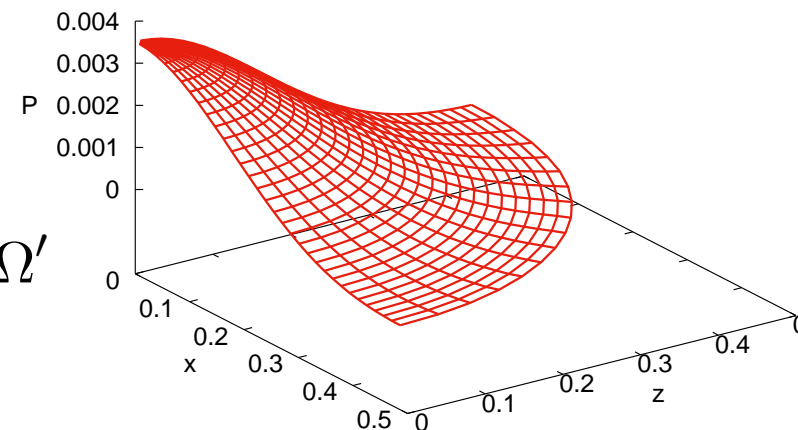
$$\mathcal{F}(\Omega) \equiv \int_{\Omega_c}^{\Omega} F(\Omega') d\Omega'$$

➔ $\nabla_{\mu} (H(p) - \log(u^t) + \mathcal{F}(\Omega)) = 0.$

Metric Component



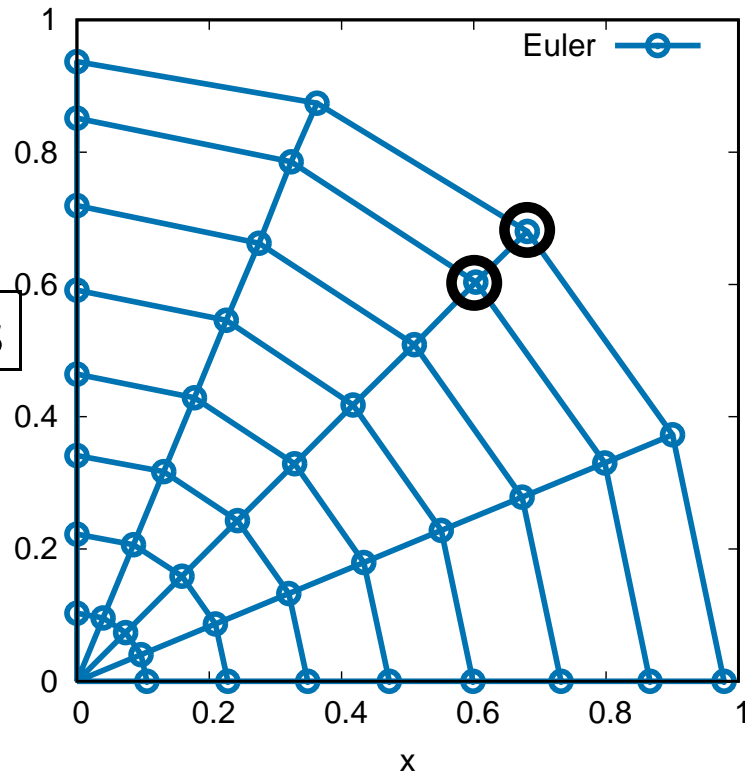
Pressure



One of the Einstein eqs.

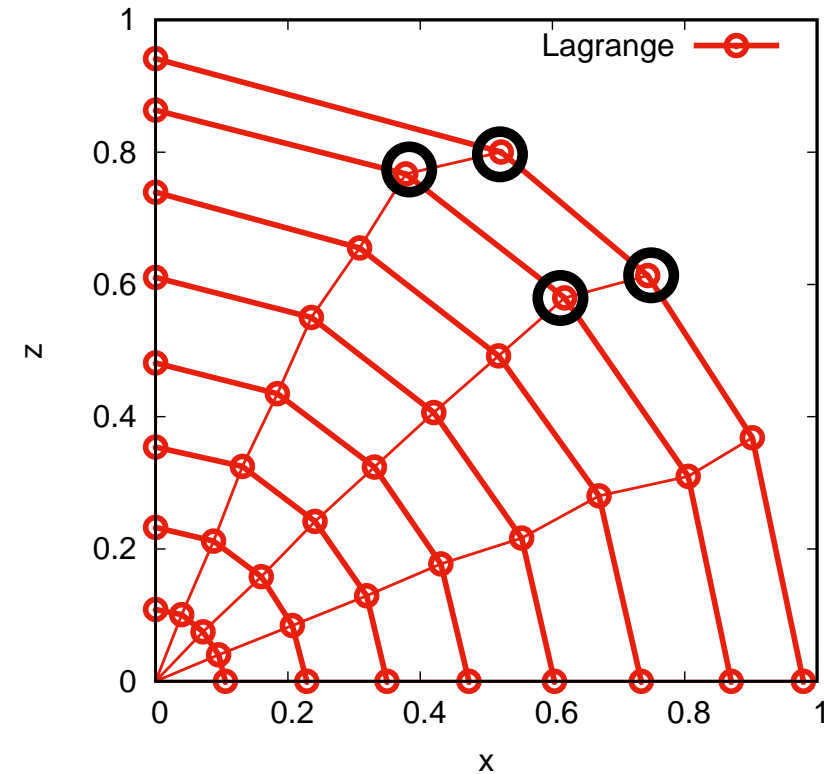
$$\begin{aligned}
 0 &\equiv G_{tt} - \frac{8\pi G}{c^4} T_{tt} \\
 &= \frac{A^2 \omega^2 \sin^2 \theta}{NB^2} \left(r^2 \frac{\partial^2 N}{\partial r^2} + \frac{\partial^2 N}{\partial \theta^2} \right) + \left(\frac{N^2 A}{r^2} - \frac{A \omega^2 \sin^2 \theta}{B^2} \right) \left(r^2 \frac{\partial^2 A}{\partial r^2} + \frac{\partial^2 A}{\partial \theta^2} \right) \\
 &\quad + \frac{N^2 A^2}{r^2 B^2} \left(r^2 \frac{\partial^2 B}{\partial r^2} + \frac{\partial^2 B}{\partial \theta^2} \right) - \frac{A^2 \omega \sin^2 \theta}{B^2} \left(r^2 \frac{\partial^2 \omega}{\partial r^2} + \frac{\partial^2 \omega}{\partial \theta^2} \right) \\
 &\quad + \frac{A^2 \omega \sin^2 \theta}{NB^2} \left(r^2 \frac{\partial N}{\partial r} \frac{\partial \omega}{\partial r} + \frac{\partial N}{\partial \theta} \frac{\partial \omega}{\partial \theta} \right) + \frac{3A^2 \omega \sin^2 \theta}{B^3} \left(r^2 \frac{\partial B}{\partial r} \frac{\partial \omega}{\partial r} + \frac{\partial B}{\partial \theta} \frac{\partial \omega}{\partial \theta} \right) \\
 &\quad - \left(\frac{N^2}{r^2} - \frac{\omega^2 \sin^2 \theta}{B^2} \right) \left\{ r^2 \left(\frac{\partial A}{\partial r} \right)^2 + \left(\frac{\partial A}{\partial \theta} \right)^2 \right\} - \frac{2N^2 A^2}{r^2 B^2} \left\{ r^2 \left(\frac{\partial B}{\partial r} \right)^2 + \left(\frac{\partial B}{\partial \theta} \right)^2 \right\} \\
 &\quad - \left(\frac{A^2 \sin^2 \theta}{4B^2} + \frac{3r^2 \omega^2 \sin^4 \theta}{4N^2 B^4} \right) \left\{ r^2 \left(\frac{\partial \omega}{\partial r} \right)^2 + \left(\frac{\partial \omega}{\partial \theta} \right)^2 \right\} \\
 &\quad + \frac{r A^2 \omega^2 \sin^2 \theta}{NB^2} \frac{\partial N}{\partial r} + \left(\frac{N^2 A}{r} - \frac{r A \omega^2 \sin^2 \theta}{B^2} \right) \frac{\partial A}{\partial r} + \frac{2N^2 A^2}{r B} \frac{\partial B}{\partial r} \\
 &\quad - \frac{4r A^2 \omega \sin^2 \theta}{B^2} \frac{\partial \omega}{\partial r} + \frac{2N^2 A^2 \cos \theta}{r^2 B \sin \theta} \frac{\partial B}{\partial \theta} - \frac{3A^2 \omega \sin \theta \cos \theta}{N^2} \frac{\partial \omega}{\partial \theta} \\
 &\quad - \frac{8\pi G}{c^4} \left[(\Omega - \omega)^2 \varepsilon \omega^2 r^4 \sin^4 \theta + (\Omega^2 P + 2\Omega \omega \varepsilon - 2\omega^2 \varepsilon) N^2 B^2 r^2 \sin^2 \theta + \varepsilon N^4 B^4 \right]
 \end{aligned}$$

Euler and Lagrange formulations



$$\rho_{i,j} \equiv \frac{\Delta m_{i,j}}{\Delta V_{i,j}}$$

$\Delta m, \Delta j, \Delta s, Y_e$
attached to
a fluid element



Discretization

$$\frac{\partial P}{\partial r} \sim \frac{P_{i,j} - P_{i-1,j}}{r_{i,j} - r_{i-1,j}}$$

PDEs

→ Nonlinear system eqs.

Standard solver for them like Newton-Raphson method does not work.

(from 1970s...)

Newton-Raphson method

To find the roots of single-variable function:

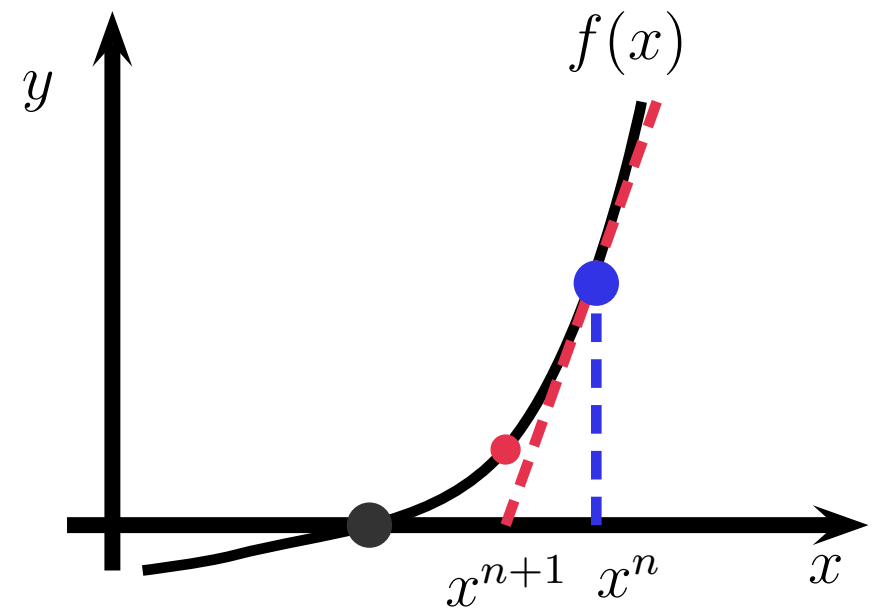
$$f(x) = 0.$$

Taylor expansion
around the solution $x^* = x + \Delta x$:

$$f(x + \Delta x) = f(x) + f' \Delta x + \mathcal{O}(\Delta x^2).$$

Since $f(x + \Delta x) = 0$, we define the map:

$$\Delta x = -\frac{f(x)}{f'},$$
$$x^{n+1} \equiv x + \Delta x = x^n - \frac{f}{f'}.$$



Linearization:

$$y - f(x^n) = f'(x^n)(x - x^n)$$

New approximated solution is expected to be closer to the solution. A series of x^n converges to the solution.

System of nonlinear eqs.(NR and W4)

(ex) Cross-section
between circle and curves

$$F_1(x, y) := x^2 + y^2 - 4 = 0,$$

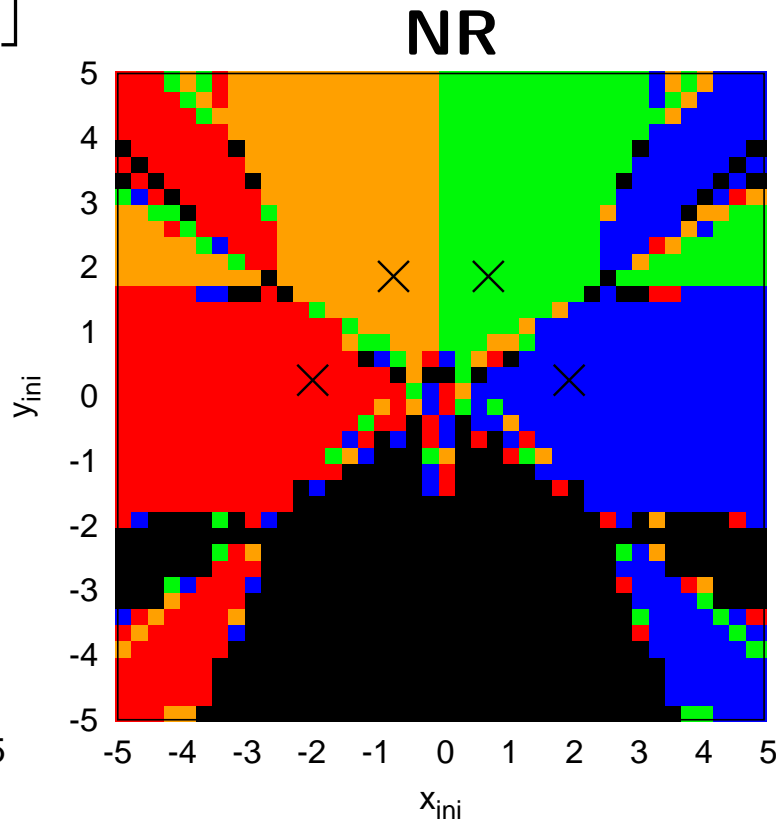
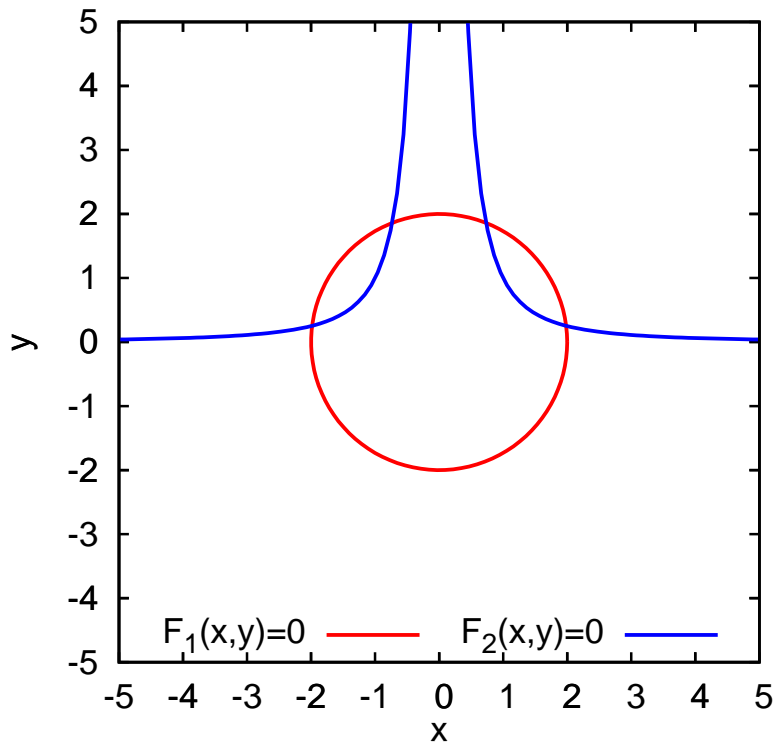
$$F_2(x, y) := x^2 y - 1 = 0.$$

$$J = \begin{bmatrix} \frac{\partial F_1}{\partial x} & \frac{\partial F_1}{\partial y} \\ \frac{\partial F_2}{\partial x} & \frac{\partial F_2}{\partial y} \end{bmatrix}$$

Newton-Raphson method

$$\left(\frac{\mathbf{x}^{n+1} - \mathbf{x}^n}{\Delta\tau} = -J^{-1} \mathbf{F} \right)$$

$$\dot{\mathbf{x}} = -J^{-1} \mathbf{F}$$



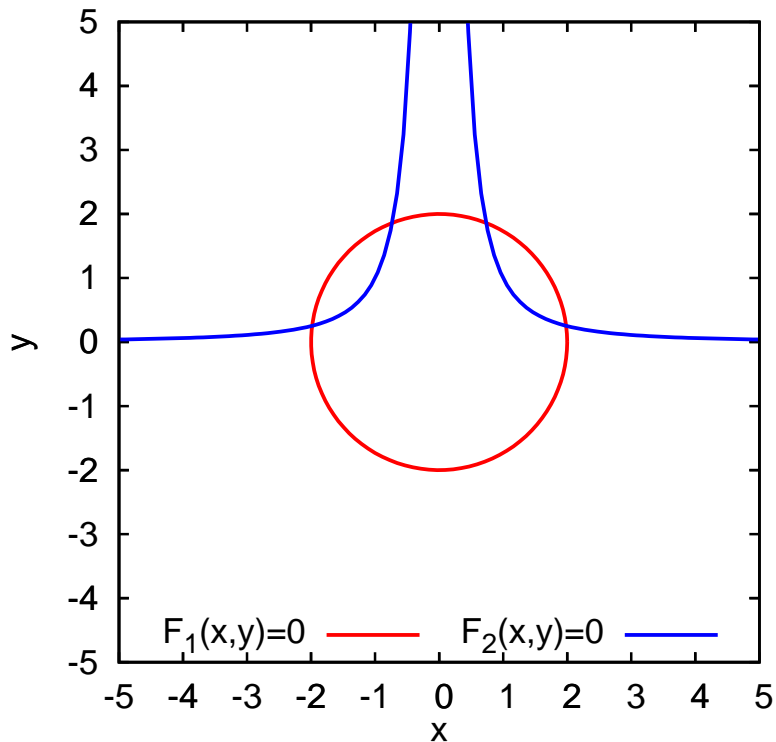
System of nonlinear eqs.(NR and W4)

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Newton-Raphson method

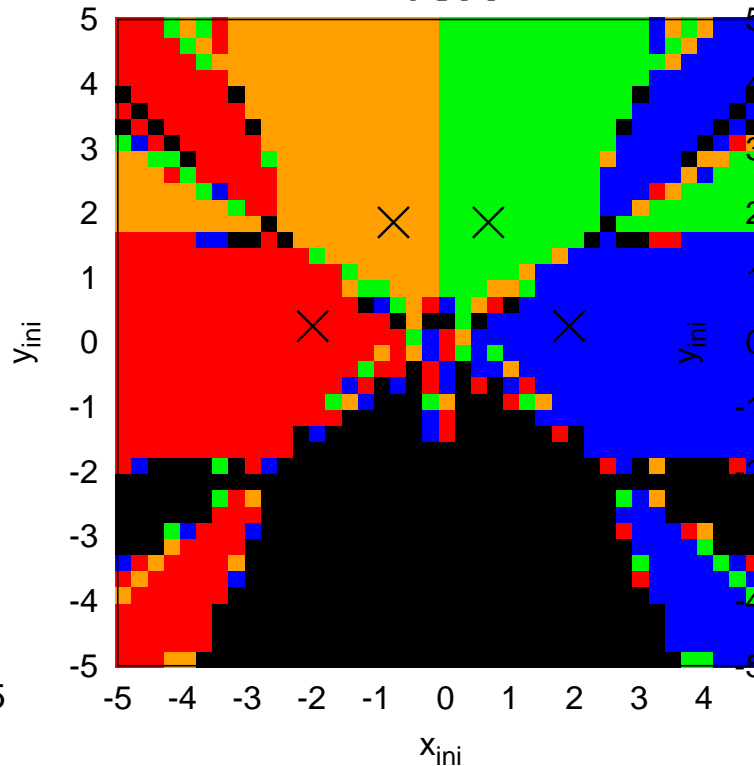
$$\left(\frac{\mathbf{x}^{n+1} - \mathbf{x}^n}{\Delta\tau} = -J^{-1} \mathbf{F} \right) \quad \dot{\mathbf{x}} = -J^{-1} \mathbf{F}$$

W4 method (Okawa *et al.* '18) ($Y = X^{-1} J^{-1}$)

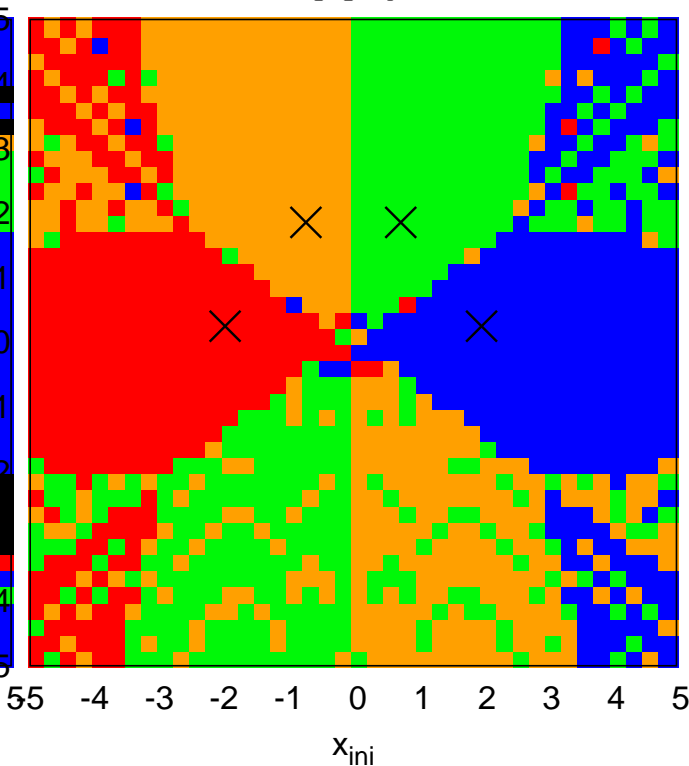
$$\dot{\mathbf{x}} = X \mathbf{p},$$

$$\dot{\mathbf{p}} = -2\mathbf{p} - Y \mathbf{F}.$$

NR



W4



Cooling of Rotating Star

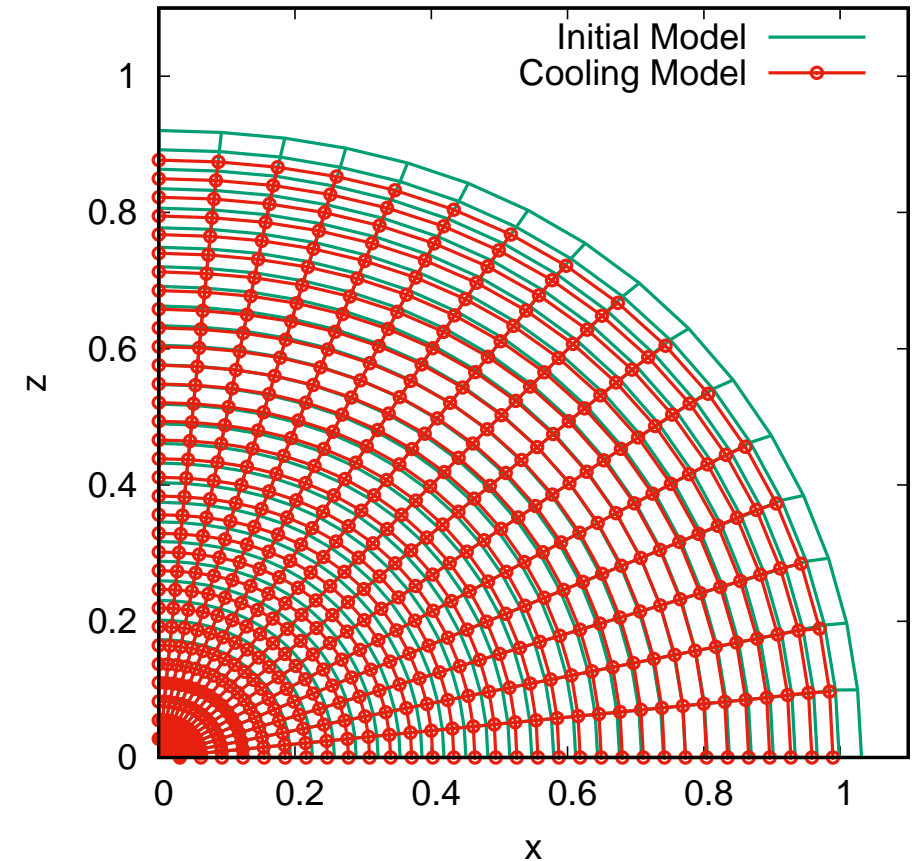
Initial model (Green curves)

- EOS: Polytrope
 $P = K_0 \rho^2$
- Rotation Law: Differential
(Constant- j law)

Evolution model (Red curves)

- Δm and Δj are Fixed.
- Temperature variation: $P = K(t) \rho^2$
- Specific angular momentum: $j_\varphi = hu_\varphi$ (Axisymmetry)

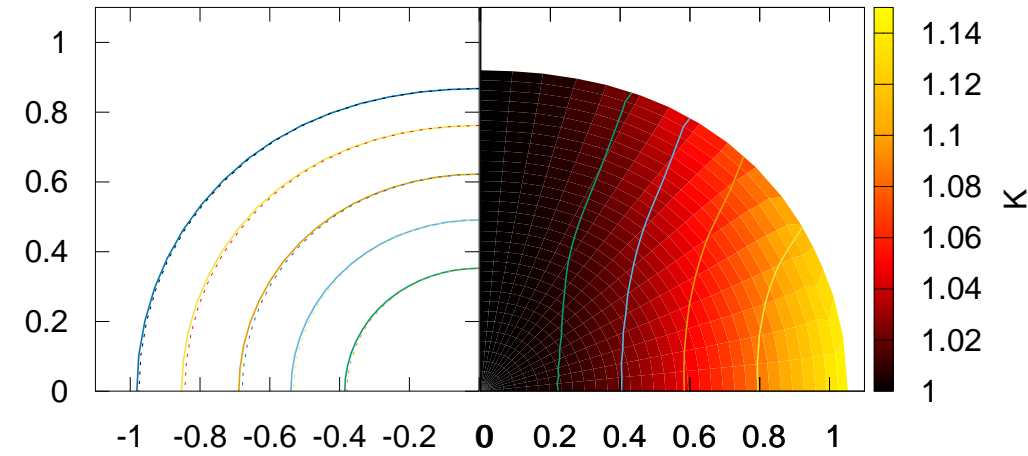
Cf. $l_\varphi \equiv -j_\varphi/j_t$:
$$l_\varphi = \frac{(\Omega - \omega)r^2 \sin^2 \theta}{N^2 + \omega(\Omega - \omega)r^2 \sin^2 \theta}$$



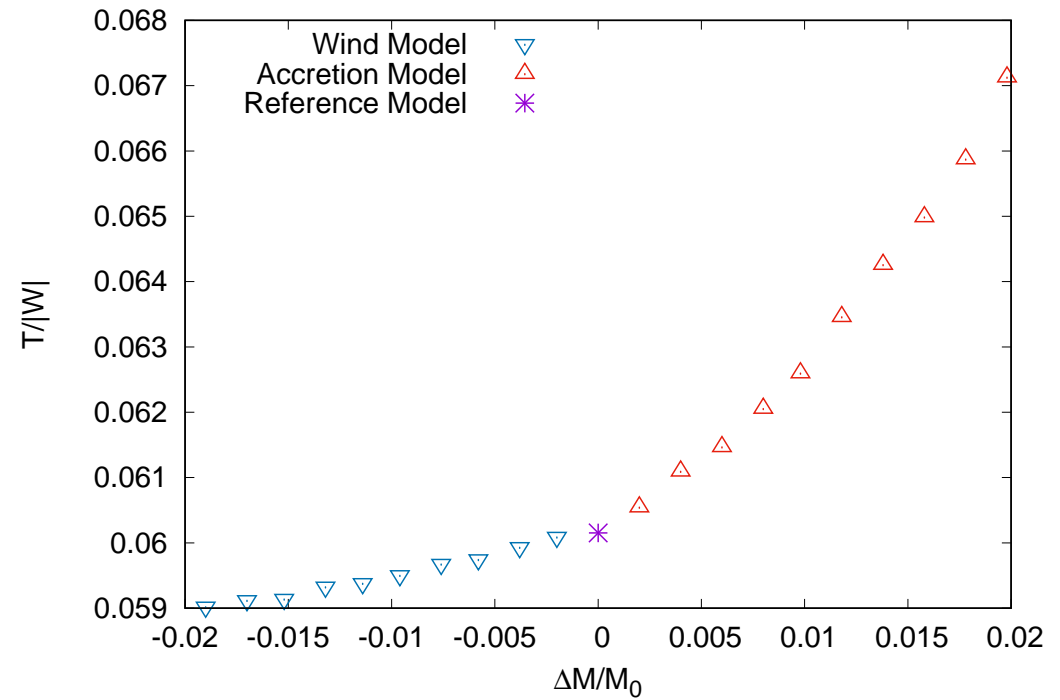
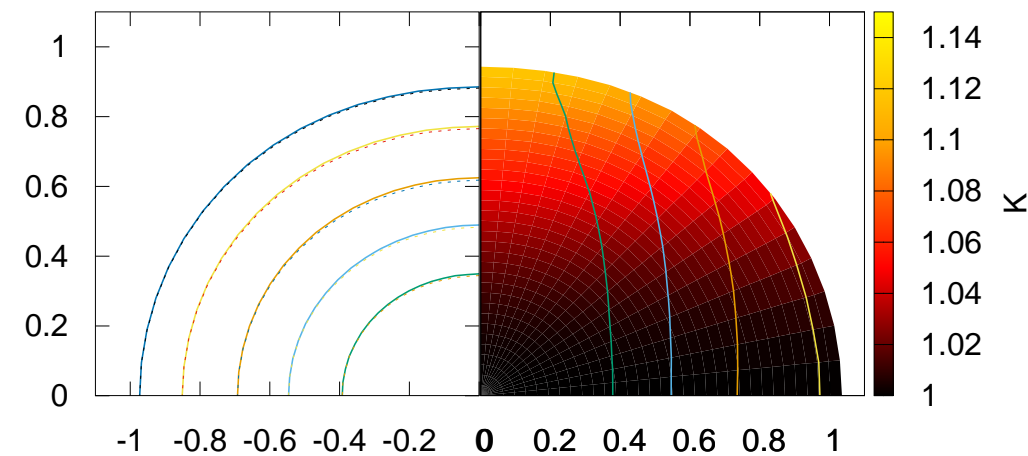
Various models for Rotating Stars in GR

Any EOS: (ex) $P(\rho, s(r, \theta)) = K(r, \theta)\rho^2$

Accretion or Mass-loss



$$K(r, \theta) = K_0 \left(1 + \epsilon \frac{r^2}{R_e^2} \sin^2 \theta \right)$$



Lagrange Evolution!!

Summary

- Equilibrium of rotating stars in GR
= Solving the system of nonlinear eqs.
 - PDEs: system of nonlinear eqs.
 - Newton-Raphson method needs fine-tuned initial guess.
- New method (W4) allows us to construct it!
 - Lagrange formulation for the evolution
 - Let us solve a realistic system!
(Neutrino transport, gravitational waves, EOS?)